

Maxwellian velocity distributions in Slow Time

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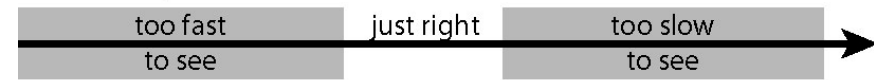
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Every observer (measuring instrument) has its own timescale:

Sally Shortwave



Laboratory scale

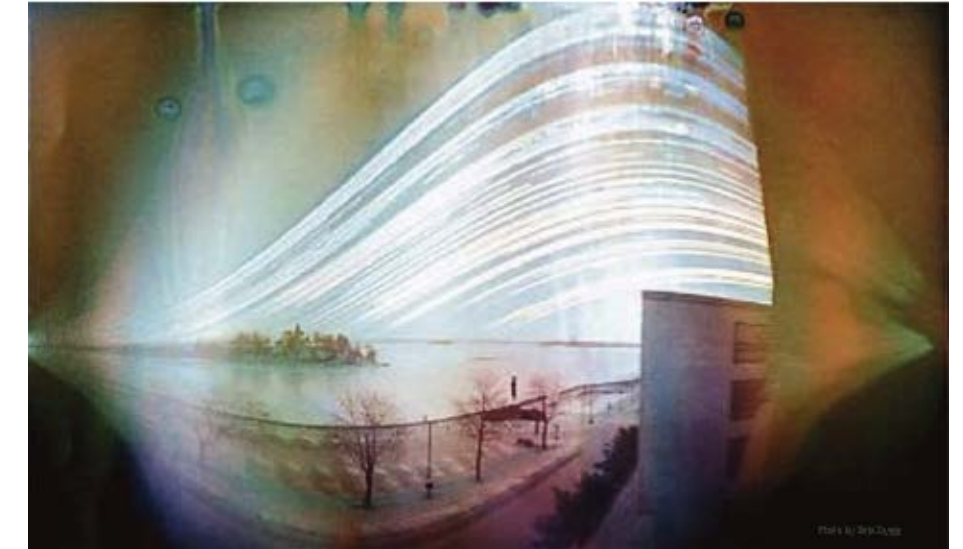


Larry Longwave



We do not see the jitter of molecular motion. Our observations average over that (typically Maxwellian) velocity distribution. What would we observe if we observe over a much longer timescale than our own? The pictures illustrate how some effects disappear while others become obvious, previously obscured by transients.

Sun over parking lot



6 mo

Busy campus intersection



1/125 sec



10min

Niagara river



0.4 sec



50 sec

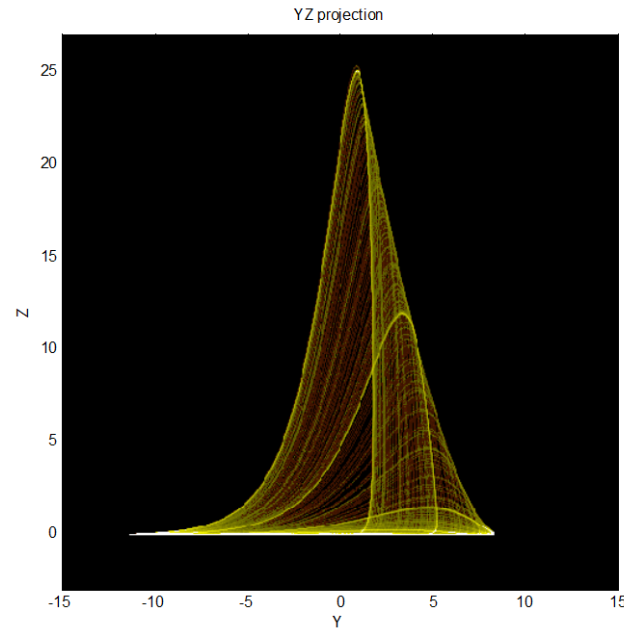
Long-time iterated maps – 10⁷ steps

Rösler equations:

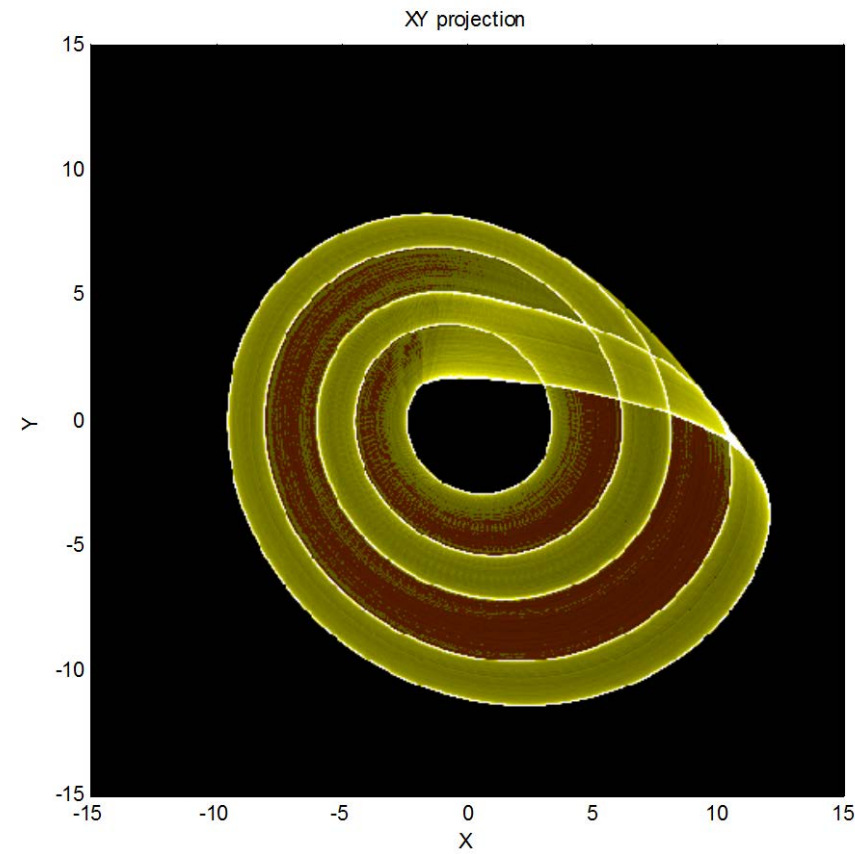
$$\dot{x} = -y - z$$

$$\dot{y} = x + 0.2y$$

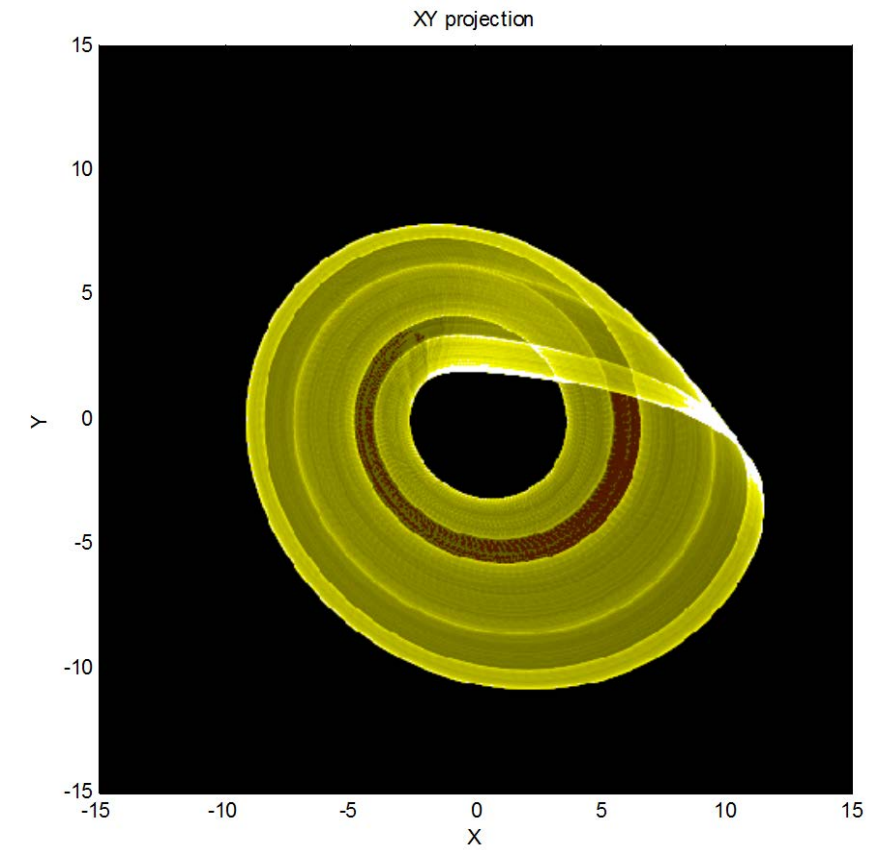
$$\dot{z} = xz - 5.2z + 0.2$$



C=5.981



C=5.982



On the long timescale we observe distributions of distributions.

Mean air velocity (wind) fluctuates and turns into a slightly increased temperature, “thermalizing” wind:

$$p(v; \theta) = \left(\frac{m}{2k\theta}\right)^{1/2} \frac{1}{\sqrt{\pi}} e^{-\frac{m}{2k\theta}v^2} \quad \theta = \frac{\sigma_u^2 m}{k} + T$$

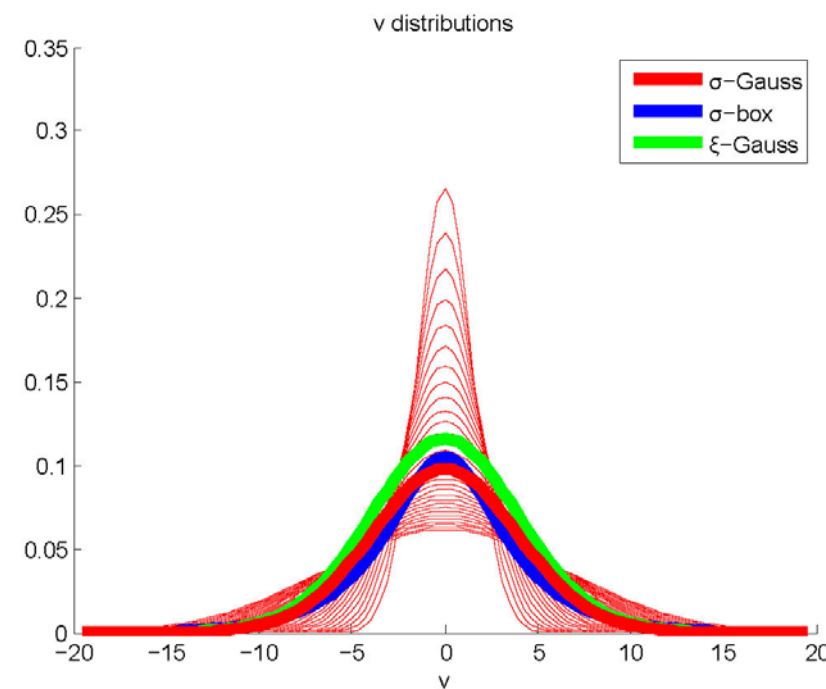
By contrast, temperature fluctuations do not result in a new temperature, a new Maxwellian distribution:

$$p(v; w, \psi_0) = \int_{-\psi_0}^{\infty} p_{v\xi} p_{\xi} d\xi = \left[\frac{1 + \operatorname{erf}\left(\frac{w^2 \psi_0}{(v^2 + w^2)^{1/2}}\right)}{1 + \operatorname{erf}(w\psi_0)} \right] \frac{w^3 \psi_0}{\sqrt{\pi}(v^2 + w^2)^{3/2}} e^{-\frac{w^2 \psi_0^2 v^2}{v^2 + w^2}} + \frac{1}{1 + \operatorname{erf}(w\psi_0)} \frac{w}{\pi(v^2 + w^2)} e^{-w^2 \psi_0^2}$$

The result is polynomial tails of degree either -3 (infinite fluctuation domain) or -2 (semi-infinite domain).

[w is the precision of the velocity precision]

Distribution of standard deviations



Convolved distribution compared to standard Maxwellian

